Measurement incompatibility and steering are necessary and sufficient for contextuality

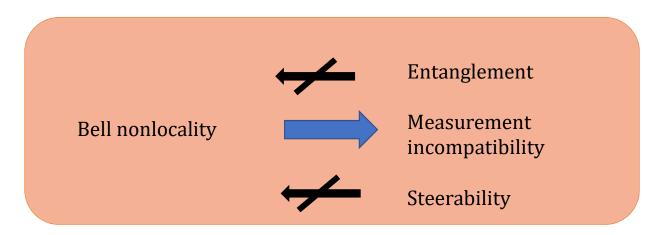
A one-to-one mapping from nonclassical entities to nonclassical correlations.

In collaboration with Roope Uola



arXiv: 1905.03614

From Bell nonlocality to nonclassical features



Joint measurability

$$M_{a|x} = \sum_{\lambda} p(a|x,\lambda)G_{\lambda}$$

<u>Unsteerability</u>

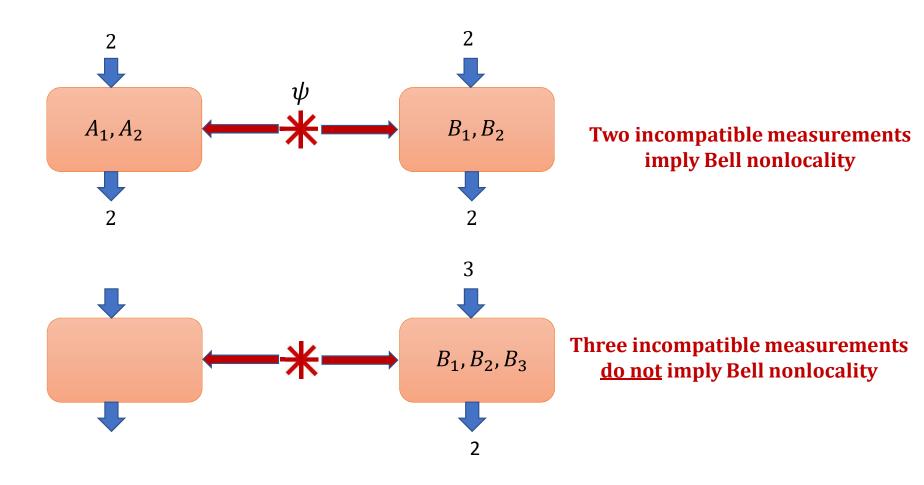
$$\sigma_{a|x} = \operatorname{Tr}_A(A_{a|x} \otimes \mathbf{1}\psi)$$

$$\sigma_{a|x} = \sum_{\lambda} p(\lambda) p(a|x, \lambda) \rho_{\lambda}$$

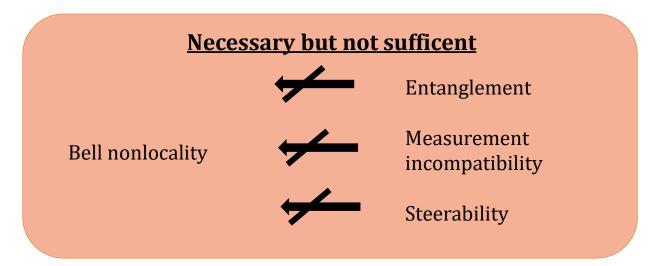
When are the nonclassical features of quantum theory also sufficient for correlations that violate a Bell inequality?

Wiseman et. al., PRL (2006). Werner PRA (1989).

Does measurement incompatibility imply Bell nonlocality?

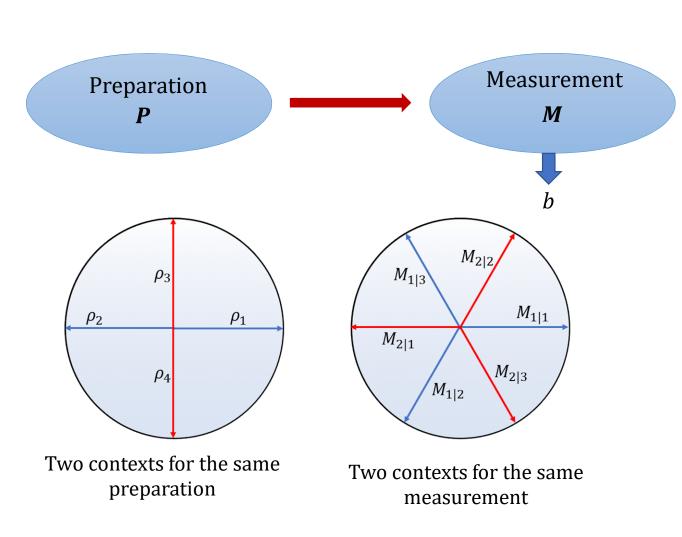


E. Bene et. al., NJP (2018). F. Hirsch et. al., PRA (2018). M. Wolf et. al., PRL (2009).



Can quantum features also be sufficient for some other form of quantum correlations?

Ontological models and contextuality



Ontological model

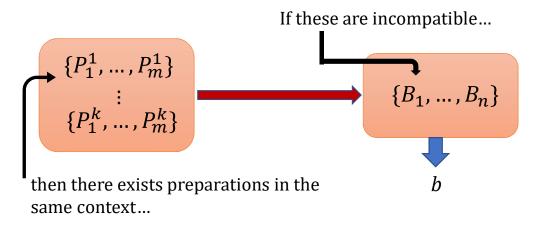
$$p(b|\mathbf{P},\mathbf{M}) = \sum_{\lambda} p(\lambda|\mathbf{P})p(b|\mathbf{M},\lambda)$$

Preparation noncontextual models are independent of the context of *P*

Measurement noncontextual models are independent of the context of *M*

Result 1

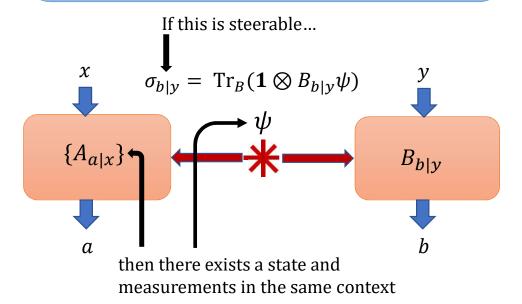
A set of measurements is jointly measurable if and only if their statistics admit a preparation noncontextual model for all states



such that the statistics is **preparation contextual**.

Result 2

An assemblage is unsteerable if and only if its statistics admits a preparation and measurement noncontextual model for all measurements



such that the statistics is contextual

Sketch of a very simple proof

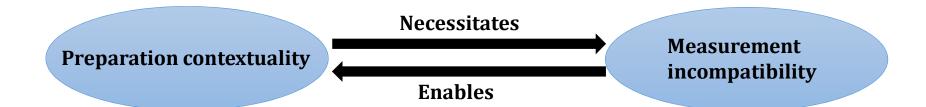
Assume that the POVM $A_{a|x}$ always produces preparation noncontextual statistics, for any state.

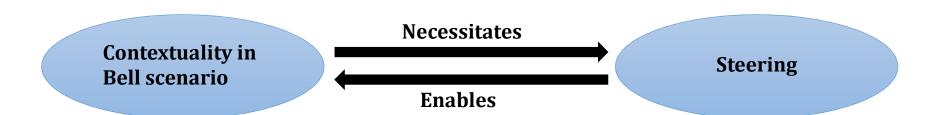
$$\forall \mathbf{P} \in \mathcal{P}_{\rho} : p(a|x, \mathbf{P}) = \sum_{\lambda} p(\lambda|\rho) p(a|x, \lambda)$$
 Convexity-preserving map from quantum states to [0,1]

Riesz representation theorem: $p(\lambda|\rho) = \text{Tr}(G_{\lambda}\rho)$ for some $0 \le G_{\lambda} \le 1$

$$\forall \mathbf{P} \in \mathcal{P}_{\rho} : p(a|x, \mathbf{P}) = \sum_{\lambda} p(a|x, \lambda) \operatorname{tr} [G_{\lambda} \rho]$$

Done





Applications

Can specific tests of contextuality ...

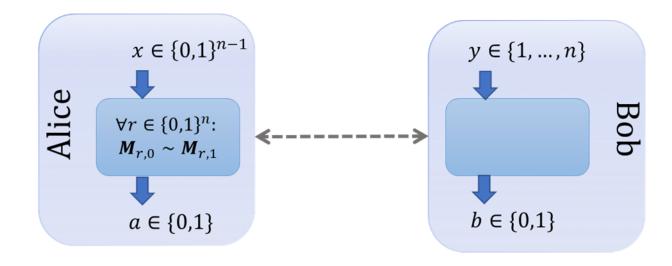
- Optimally certify steerability of interesting states?
- Certify the incompatibility of an interesting class of mesurements?

$$\sum_{x,y} p(a+b = (a,x+a)_y | x,y) \le 2^{n-2}(n+1)$$

Class of states

$$\psi = v|\psi^{-}\rangle\langle\psi^{-}| + \frac{1-v}{4}\mathbf{1}$$

$$v_2 = 0.7071$$
 $v_3 = 0.5774$ $v_4 = 0.5547$ $v_5 = 0.5422$ $v_6 = 0.5270$ $v_7 = 0.5234$.



Class of measuremets

Dichotomic qubit measurements

Numerics: 10 000 random sets of incompatible measurements for n=2,...,7

J. Bavaresco, PRA (2017)

Conclusions

 Measurement incompatibility and steering are necessary and sufficient for contextuality

 Specific tests of contextuality can optimally certify interesting classes of states and measurements.

 Certification of incompatible measurements when no Bell inequality violation is possible.

