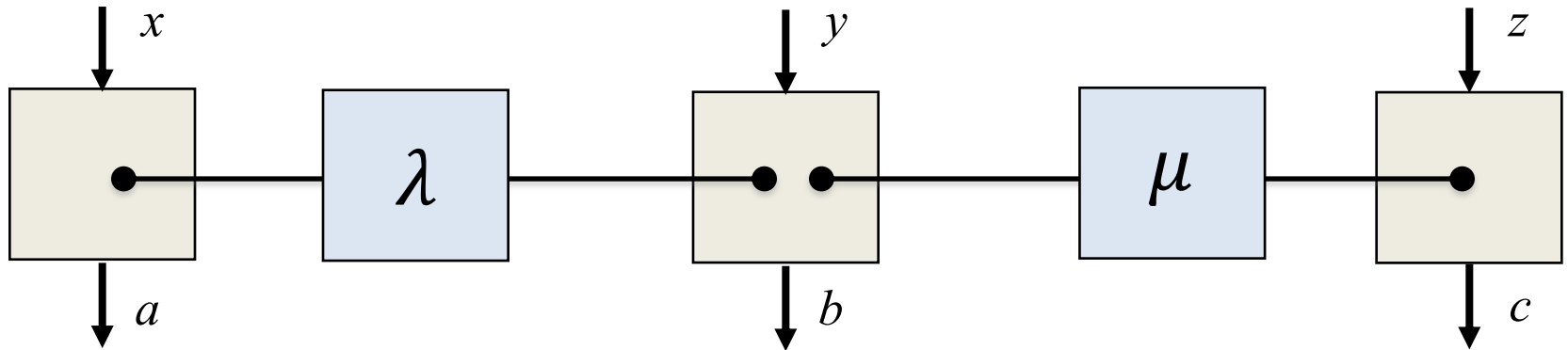


# Bounding the sets of classical and quantum correlations in networks

A. Pozas-Kerstjens et al., to appear in PRL

# Scalar extension

The method is only applicable to networks in which there are factorization constraints among observed variables, such as entanglement swapping.



$$\langle AC \rangle = \langle A \rangle \langle C \rangle$$

# Scalar extension

Example:  $S = \{1, A_0A_1, C_0C_1, \langle A_0A_1 \rangle\}$

$$\Gamma = \begin{array}{c} \mathbb{1} \\ (A_0A_1)^\dagger \\ (C_0C_1)^\dagger \\ \langle A_0A_1 \rangle^* \mathbb{1} \end{array} \begin{pmatrix} \mathbb{1} & A_0A_1 & C_0C_1 & \langle A_0A_1 \rangle \mathbb{1} \\ 1 & v_1 & v_2 & v_3 \\ & 1 & v_4 & v_5 \\ & & 1 & v_6 \\ & & & v_7 \end{pmatrix}$$

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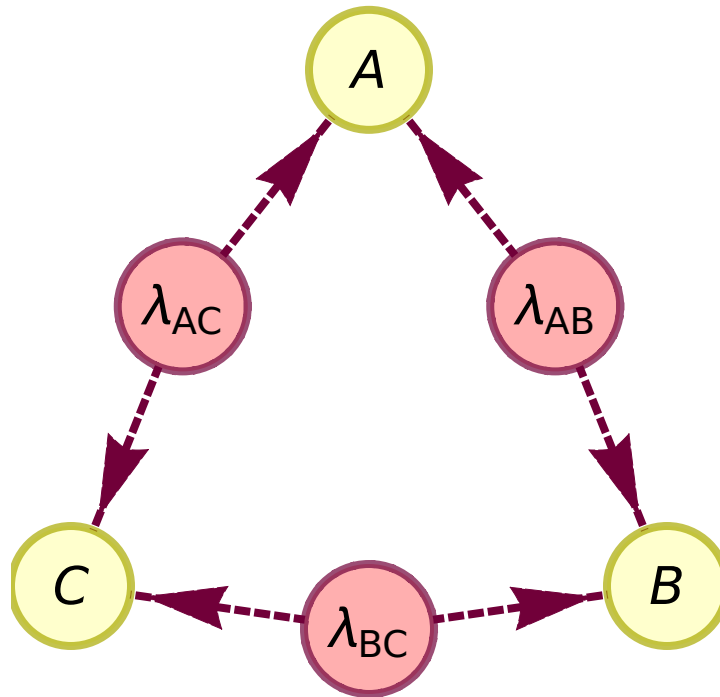
$$v_5 = v_7$$

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# Quantum inflation: a general approach to quantum causal compatibility

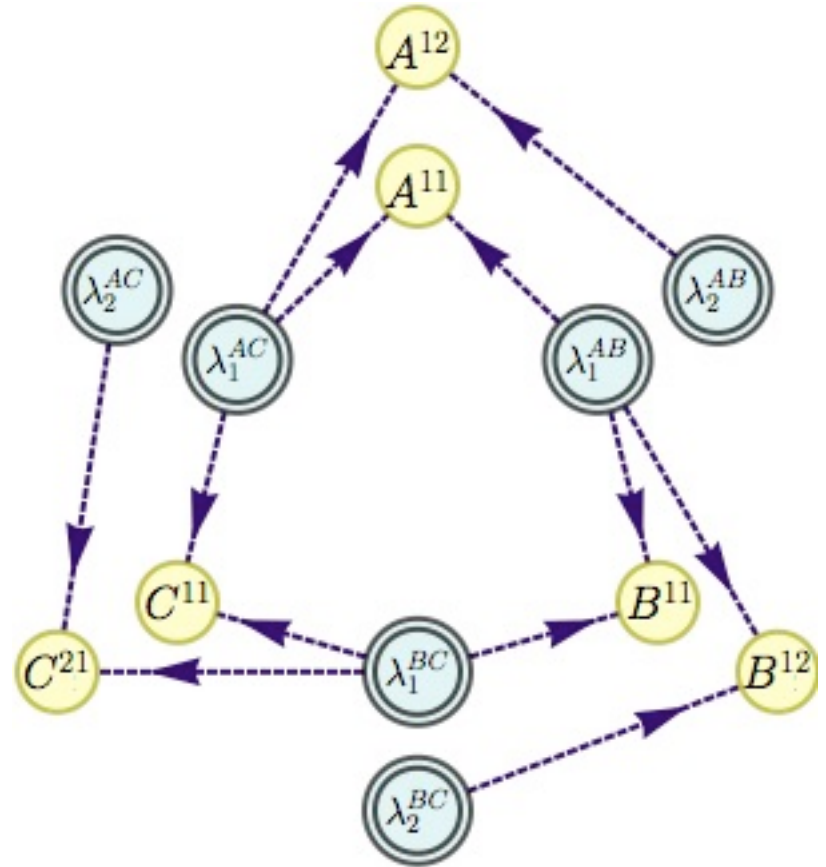
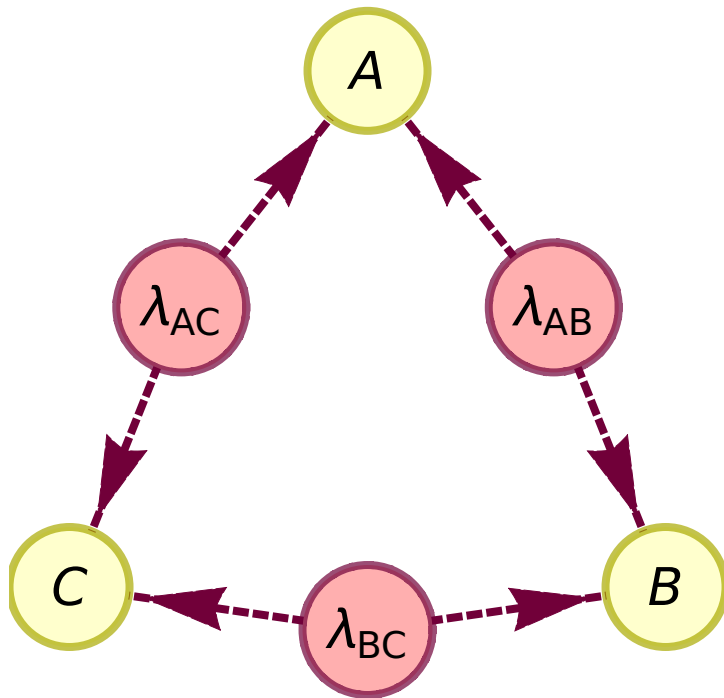
E. Wolfe et al., to appear

# Classical inflation

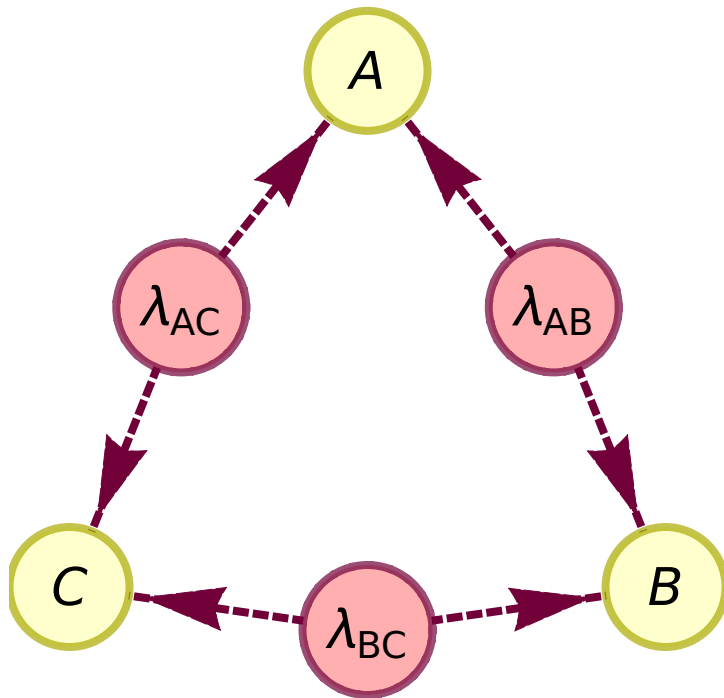




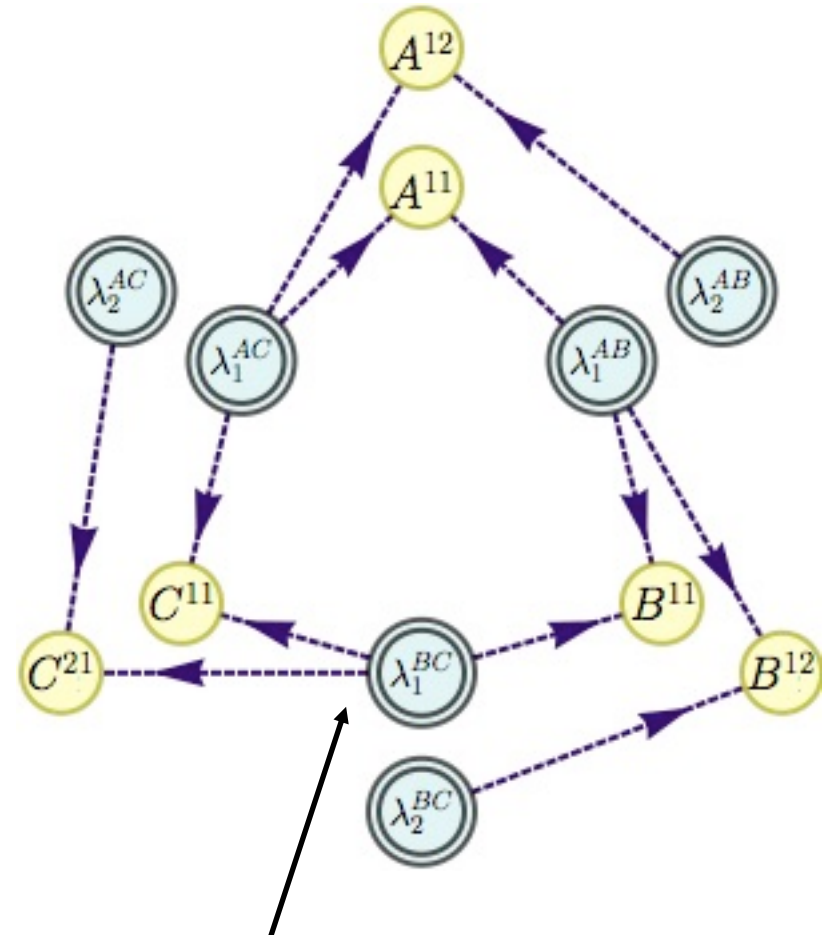
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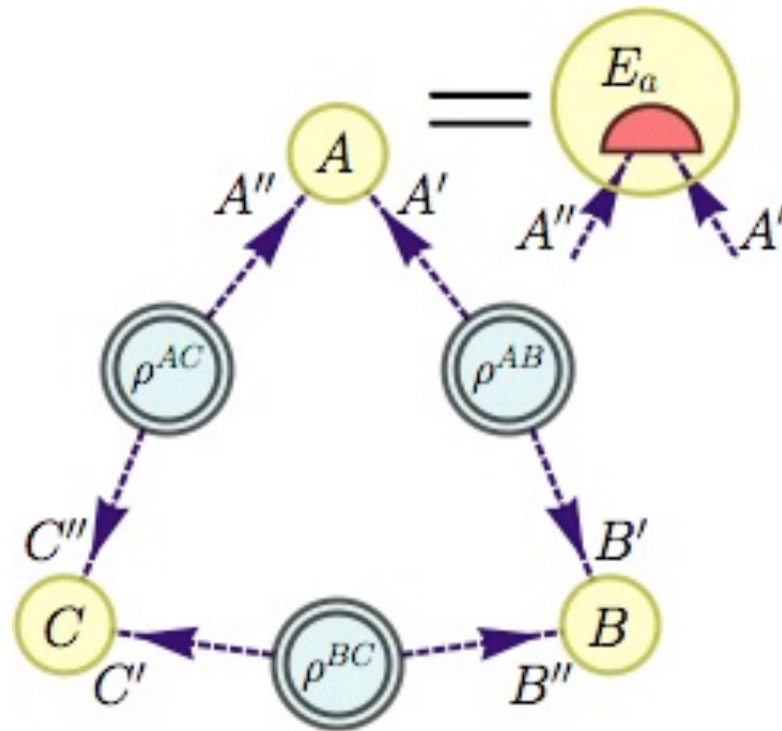


Wolfe, Spekkens & Fritz, arXiv:1609.00672

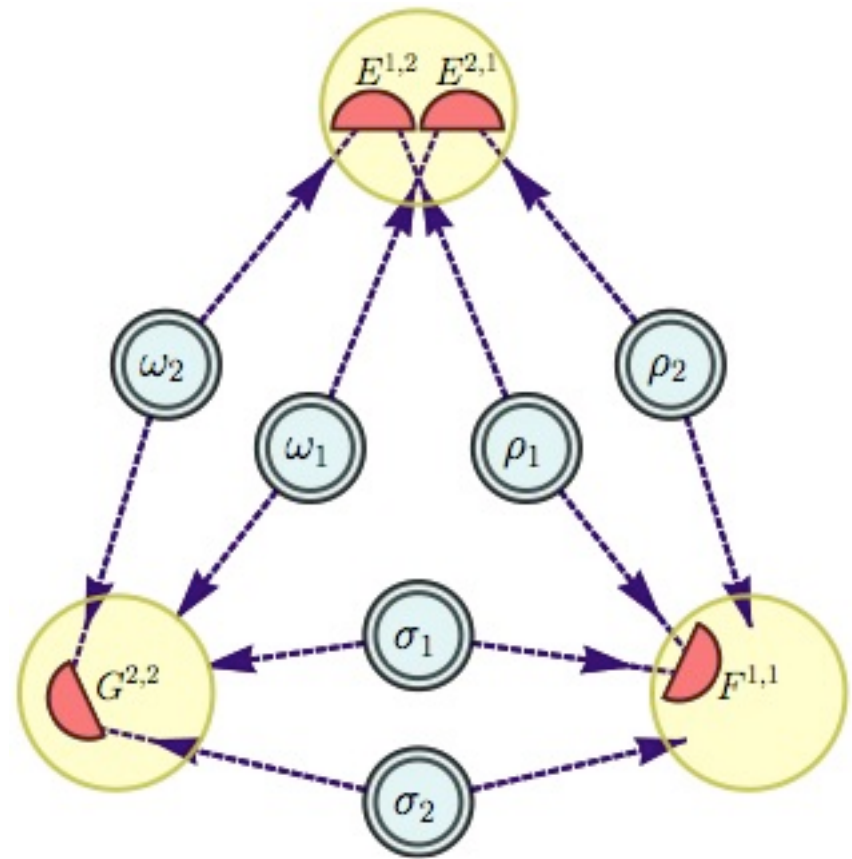
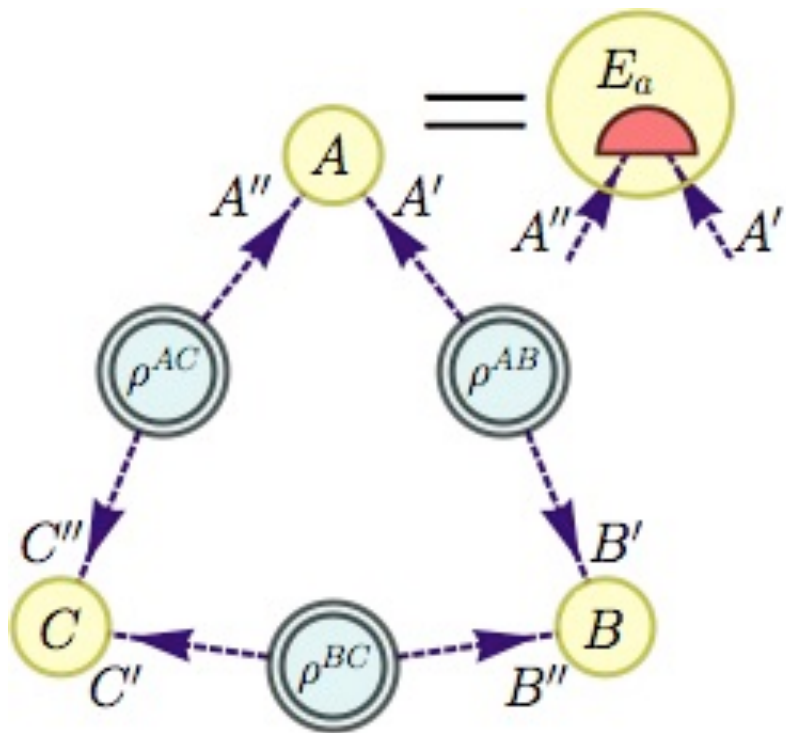


**Problem: information broadcasting!!**

# Quantum inflation



# Quantum inflation



# Quantum inflation

- Projection rules
- Commutation rules
- Symmetry under permutation of indices
- Consistency with the observed probabilities

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- Symmetry under permutation of indices
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All these conditions can be imposed with non-commutative polynomial optimization.

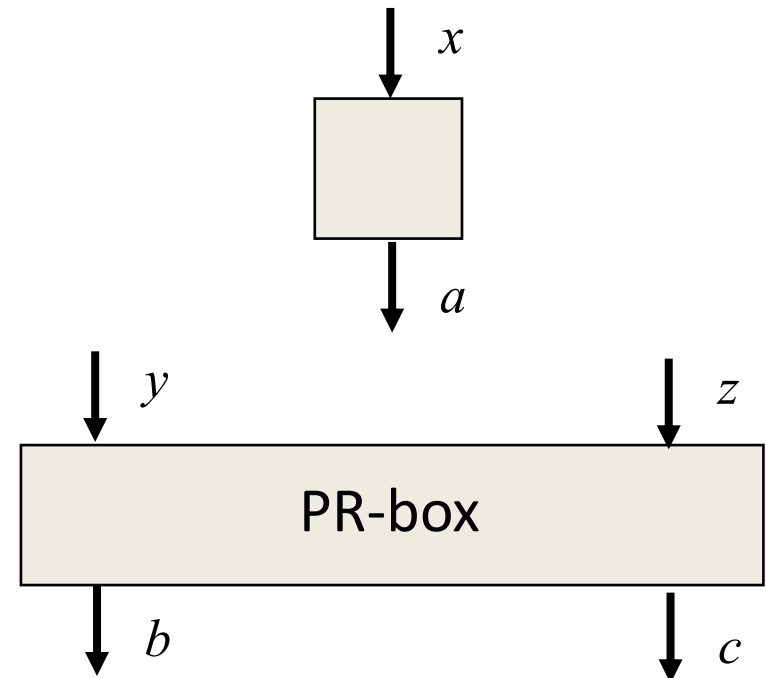
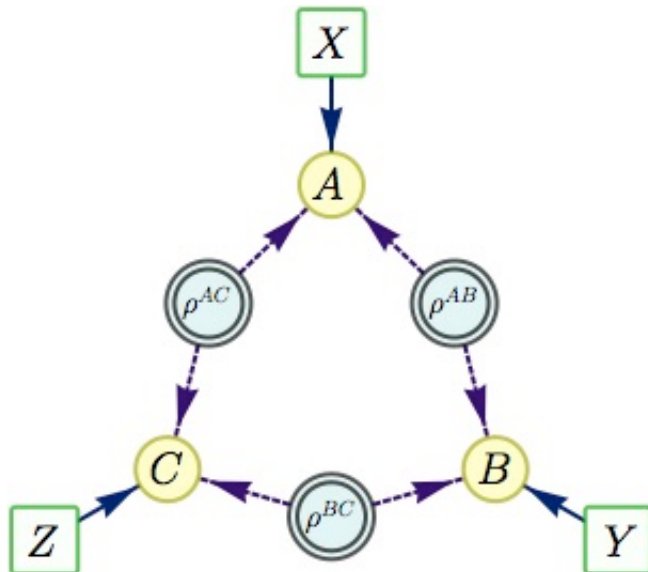
$$\begin{aligned} p^* &= \min_{(\mathcal{H}, X, \rho)} \langle p(X) \rangle_\rho \\ \text{s.t. } & q_i(X) \succeq 0 \quad \forall i = 1 \dots m_q \end{aligned}$$

# Results

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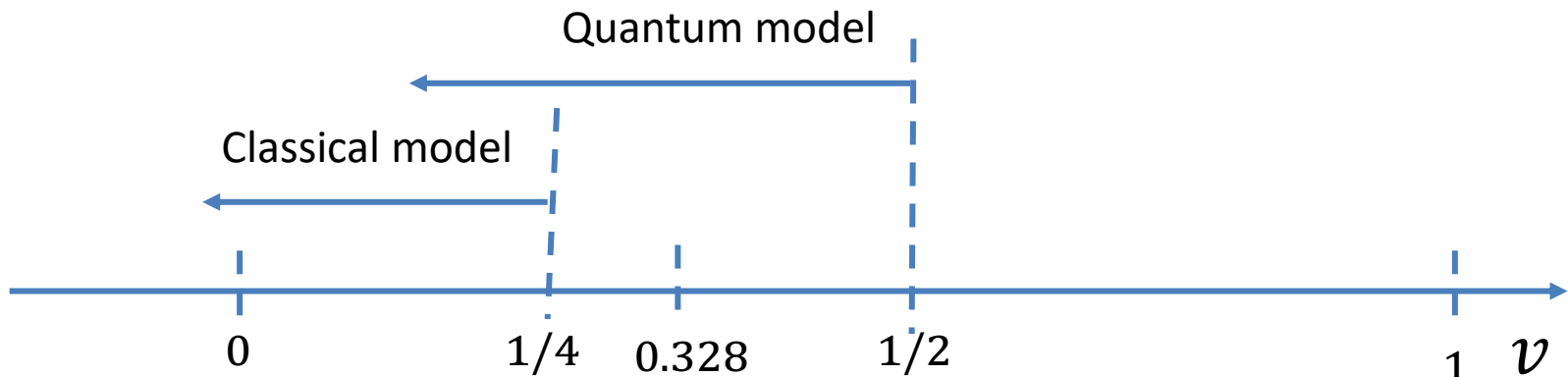


# Results

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- Inflation + scalar: known quantum (both) and classical (scalar) violations for entanglement swapping..

$$P_{2\text{PR}}(a, b, c | x, y, z) := [1 + (-1)^{a+b+c+xy+yz}] / 8$$

$$P_{2\text{PR},v} := vP_{2\text{PR}} + (1 - v)/8$$

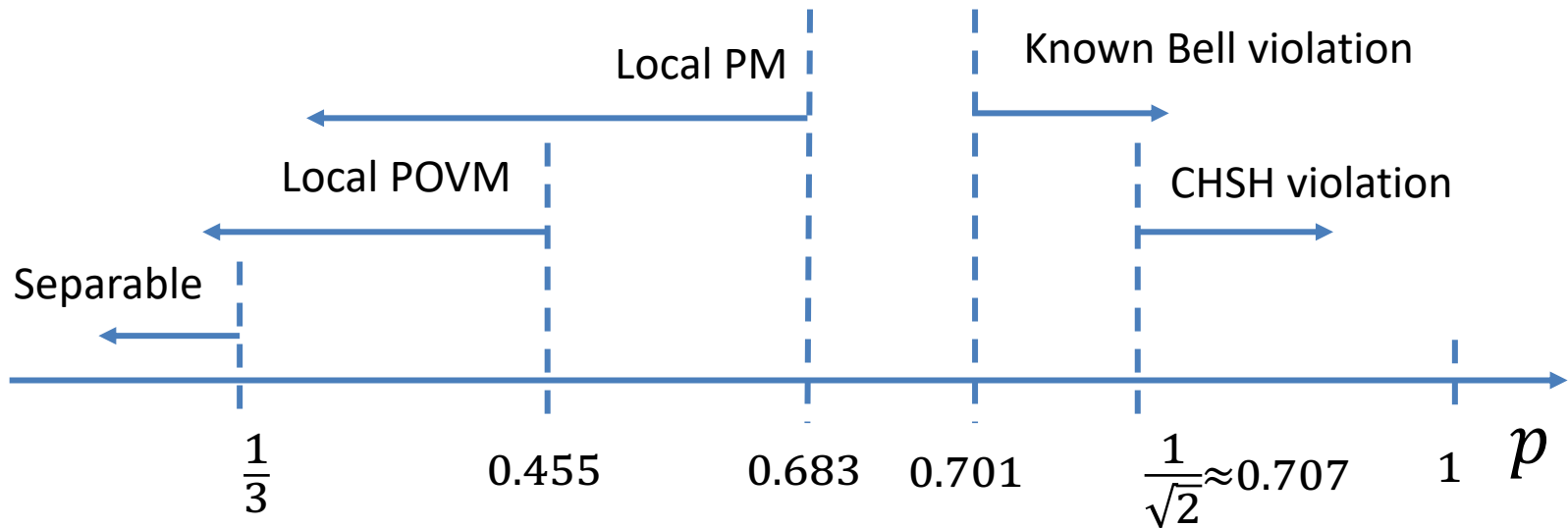


# Activation phenomena in networks

# Entanglement $\neq$ Nonlocality

Werner: there exist entangled states that do not violate any Bell inequality.

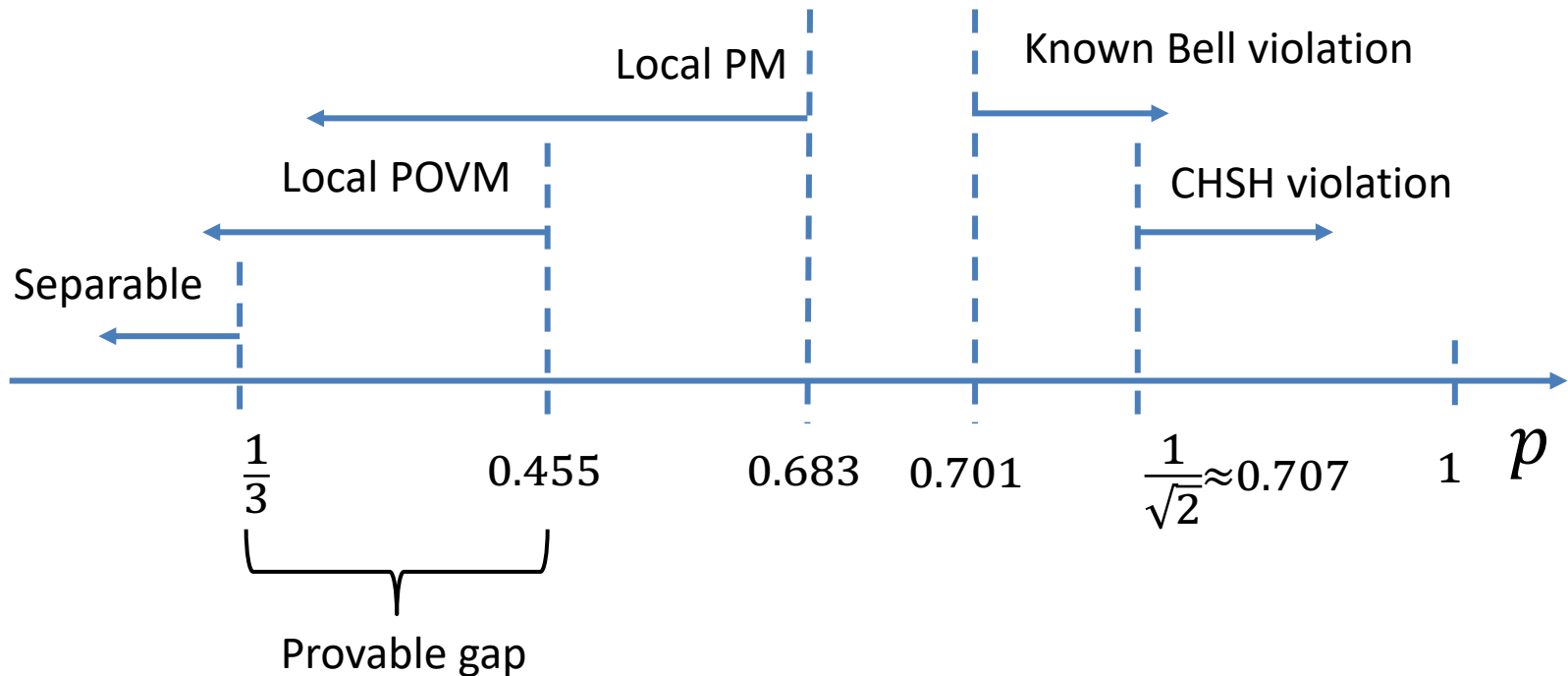
$$\rho(p) = p|\Phi\rangle\langle\Phi| + (1-p)\frac{1}{4}$$



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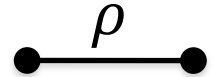
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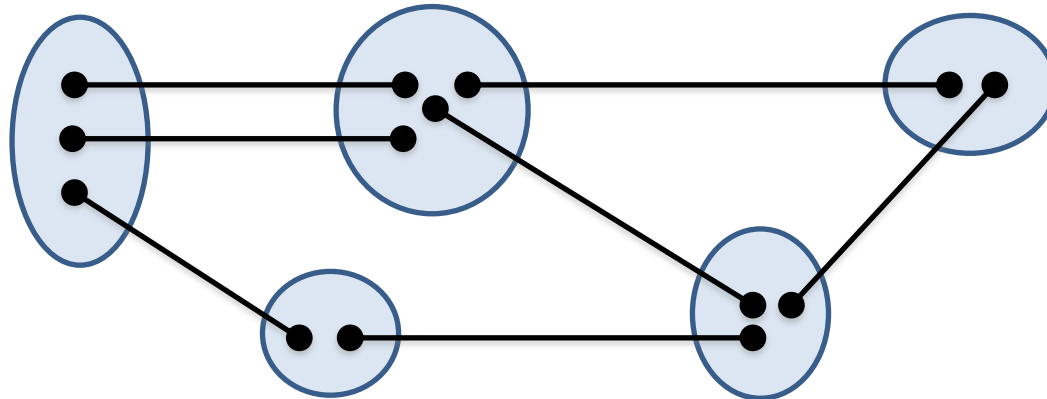
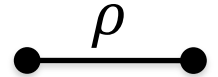
# Networks reveal non-locality

- Consider a bipartite state  $\rho$  that does not violate any Bell inequality.



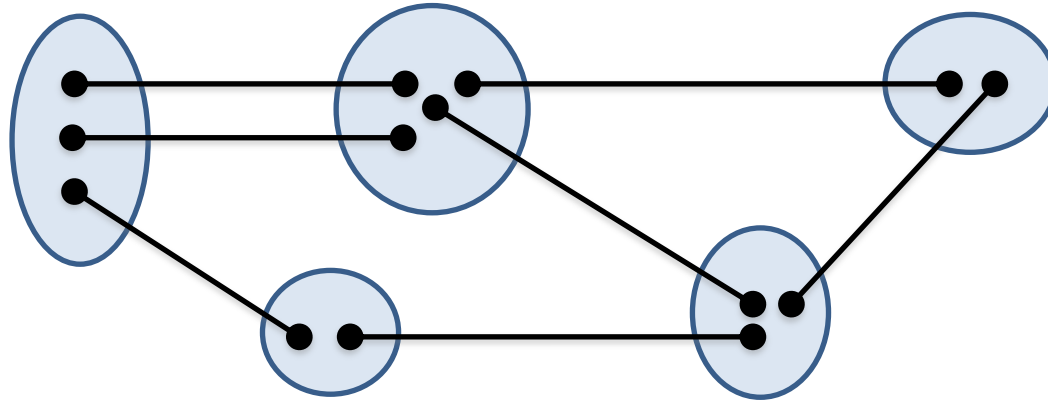
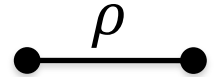
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- If the network state  $\rho_N = \rho^{\otimes N}$  violates an N-partite Bell inequality, the state  $\rho$  must be non-local.

Cavalcanti, Almeida, Scarani, Acin, Nature Comm. 2, 184 (2011)

# Networks reveal non-locality

- All bipartite states with one-way distillable entanglement are non-local.

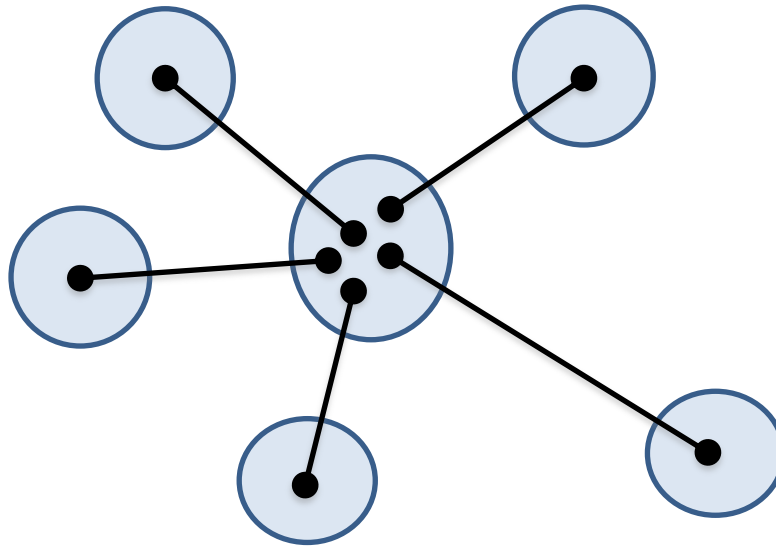


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- Non-locality of Werner state for  $p < 0.68$  in a star network with  $N > 21$ . In the limit when  $N \rightarrow \infty$ , one has  $p \rightarrow 0.64$ .



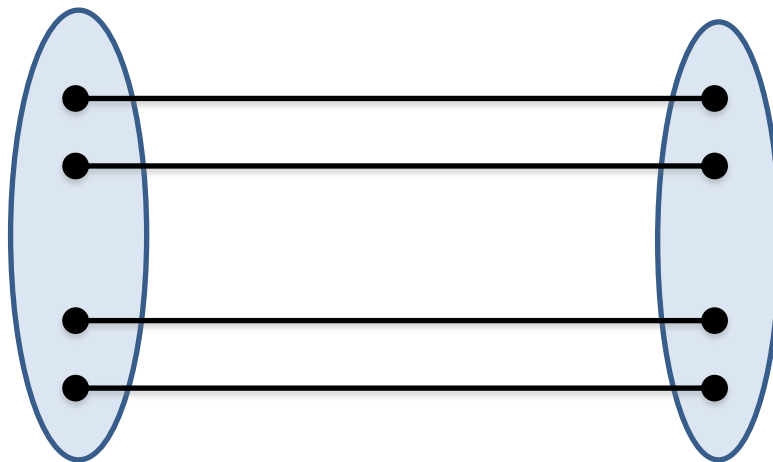
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# Activation of non-locality

- Palazuelos: there exist local entangled states  $\rho_{AB}$  such that  $\rho_{AB}^{\otimes N}$  violates a Bell inequality for large  $N$ . Palazuelos, Phys. Rev. Lett. 109, 190401 (2012)

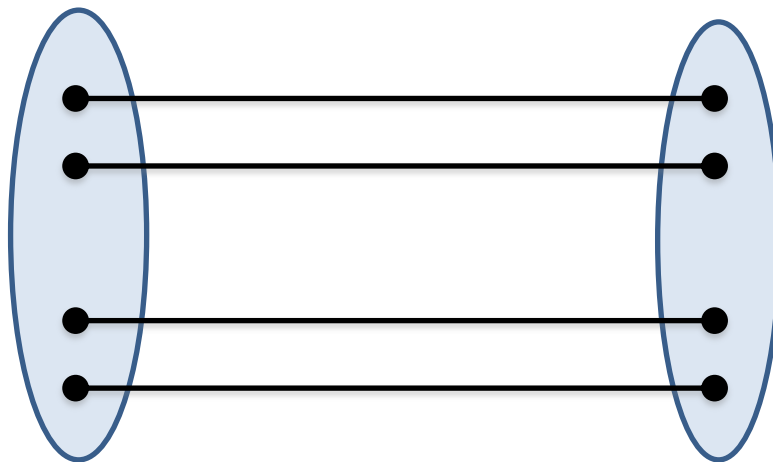
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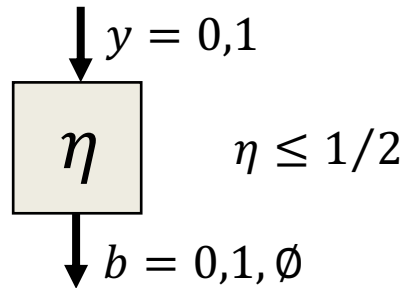
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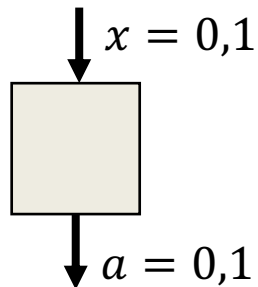
- We improved over Palazuelos' result showing that it applies to all states with singlet fidelity larger than  $1/d$ . For isotropic states, this coincides with the separability bound. [Cavalcanti, Acin, Brunner, Vertesi, PRA 87, 042104 \(2013\)](#)

# New activation phenomena

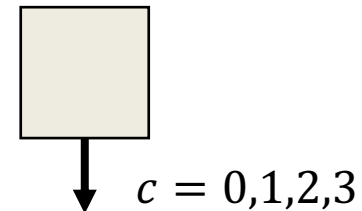
$\eta$ -CHSH Box



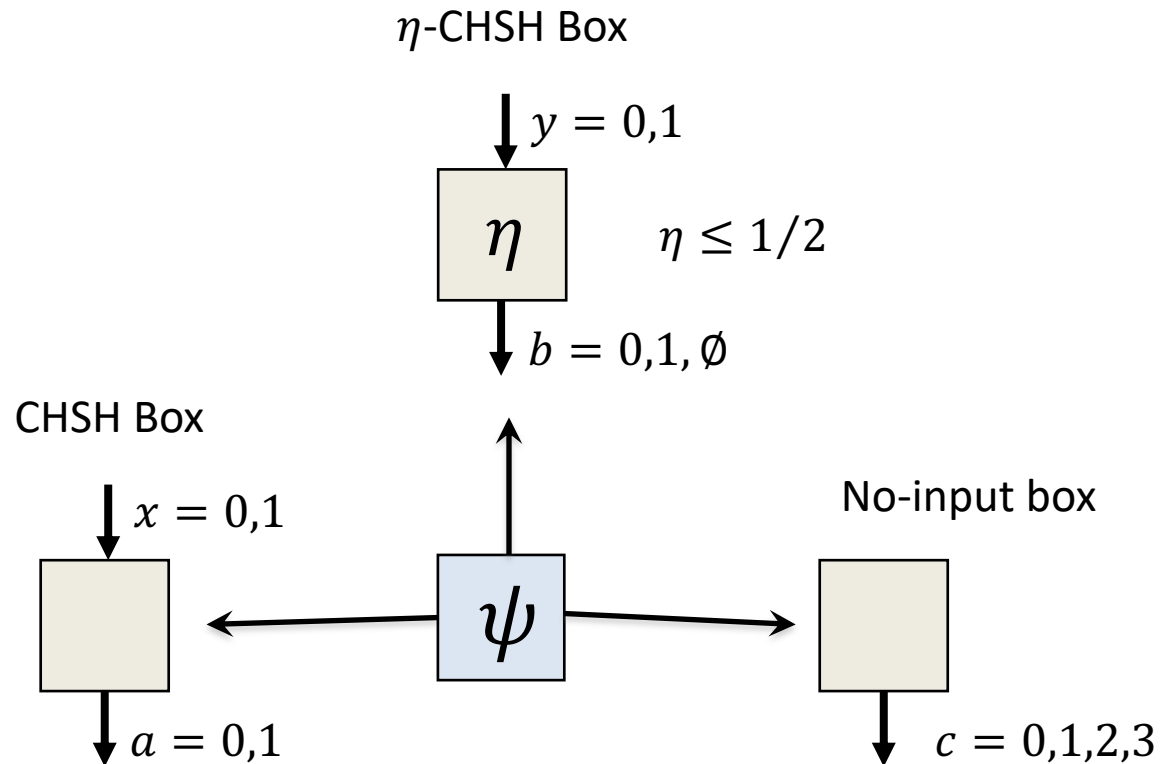
CHSH Box



No-input box

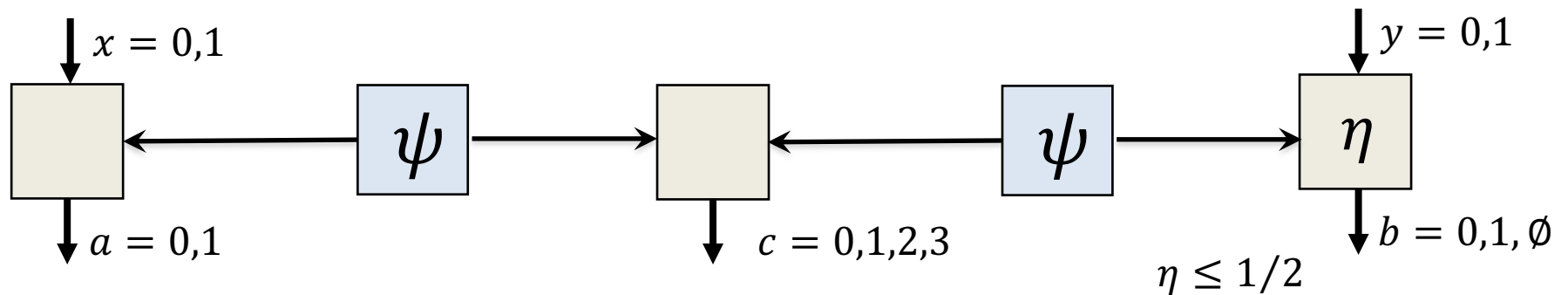


# New activation phenomena



No Bell violation is possible with these measuring devices.

# New activation phenomena



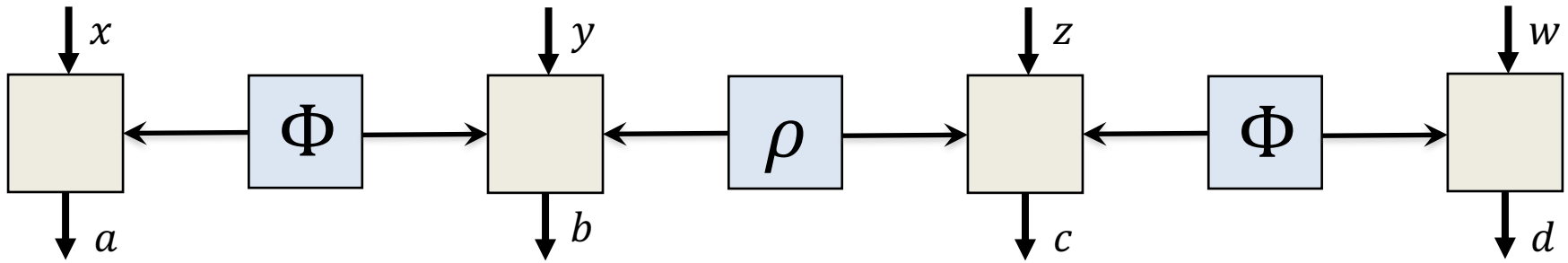
It is possible to violate bilocality with these boxes. Conjecture: valid for all  $\eta > 0$ .

To do so, we derive a method to bound the set of classical and quantum correlations in these networks through semi-definite programming.

A. Pozas-Kerstjens *et al.*, arXiv:1904.08943



# New DI entanglement detection



For any entangled state  $\rho$  one can construct a Bell inequality

$\beta_\rho = \sum W_{abcdxyzw} p(abcd|xyzw)$  such that:

- $\beta_\rho < 0$  for some  $p(abcd|xyzw) = \text{tr}(\Pi_{a|x} \otimes \Pi_{b|y} \otimes \Pi_{c|z} \otimes \Pi_{d|w} \Phi \otimes \rho \otimes \Phi)$
- $\beta_\rho \geq 0$  for all  $p(abcd|xyzw) = \text{tr}(\Pi_{a|x} \otimes \Pi_{b|y} \otimes \Pi_{c|z} \otimes \Pi_{d|w} \Phi \otimes \rho_S \otimes \Phi)$

Bowles, Supic, Cavalcanti, Acin, Phys. Rev. Lett. 121, 180503 (2018)

# Conclusions

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- Inflation method: systematic method to study quantum causality.
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- Important: small quantum networks are within reach with present technology.

# Are you interested in these topics?

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